



# Light QCD exotics at BESIII

Linjian Li (on behalf of the BESIII Collaboration) Institute of High Energy Physic, Chinese Academy of Sciences

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## QCD allows the existence of exotic hadrons

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A lot of exotic states observed experimentally, the nature of them is far from being understood

## World's Largest $\tau$ -charm Data Sets in $e^+e^-$ Annihilation

**Beijing Electron Positron Collider (BEPCII)** 



Ideal lab for light hadron physics

- Clean high statistics data samples
- Well defined initial and final states
  - Kinematic constraints
  - $I(J^{PC})$  filter
- "Gluon-rich" processes  $\Gamma(J/\psi \to \gamma G) > \Gamma(J/\psi \to \gamma H) > \Gamma(J/\psi \to \gamma M)$



### Glueballs

- Low-lying glueballs with ordinary  $J^{PC}$  (0<sup>++</sup>, 2<sup>++</sup>, 0<sup>-+</sup>)
- $\rightarrow$  mixing with  $q\bar{q}$  mesons
  - > Non- $q\bar{q}$  nature difficult to established: Cryptoexotic
    - Supernumerary states
    - Unusual pattern of production and decay
- Scalar glueball is expected to have a large production in  $J/\psi$ radiative decays:  $B(J/\psi \rightarrow \gamma G_{0^+})=3.8(9) \times 10^{-3}$  by Lattice QCD
  - Observed  $B(J/\psi \rightarrow \gamma f_0(1710))$  is  $\times 10$  larger than  $f_0(1500)$

PRD 87 092009 (2013)

PRD 92 052003 (2015)

> BESIII:  $f_0(1710)$  largely overlapped with scalar glueball

phenomenology studies of coupled channel analysis with BESIII results: PLB 816, 136227 (2021), EPJC 82, 80 (2022)





### More scalar mesons at BESIII



- $a_0(1817)$  could be the isospin one partner of the SU(3)<sub>f</sub> octet  $f_0(1800)$
- The possibility of the  $f_0(1710)$  as the scalar glueball cannot be excluded

## Light hadrons with exotic quantum numbers

- Unambiguous signature for exotics
  - Light Flavor-exotic hard to establish
  - Exotic quantum numbers forbidden for simple  $q\bar{q}$  system
  - Efforts concentrate on Spin-exotic

• 
$$J^{PC} = 0^{--}$$
, even<sup>+-</sup>, odd<sup>-+</sup>(1<sup>-+</sup>)



## Hybrids

- Low-lying hybrids can have **exotic quantum numbers**
- The exotic  $J^{PC} = 1^{-+}$  nonet of hybrids is predicted to be the lightest
  - Only isovector candidates  $\pi_1(1400), \pi_1(1600), \pi_1(2015)$  observed yet
    - \*  $\pi_1(1400)$ ,  $\pi_1(1600)$  can be explained as one resonance with recent coupled channel analysis
- **Isoscalar**  $1^{-+}$  is critical to establish the hybrid nonet
  - Can be produced in the gluon-rich  $J/\psi$  radiative decays
  - Can decay to  $\eta \eta'$  in P-wave

PRD 83,014021 (2011) PRD 83,014006 (2011) Eur.Phys.J.Plus 135, 945(2020)

#### Lattice QCD Predictions:



### **Observation of An Exotic Isoscalar State** $\eta_1(1855)$ (1<sup>-+</sup>) in $J/\psi \rightarrow \gamma \eta \eta'$

• Partial wave analysis of  $J/\psi \rightarrow \gamma \eta \eta'$ 

**Quasi two-body decay amplitudes** in the sequential decay processes  $J/\psi \rightarrow \gamma X, X \rightarrow \eta \eta'$  and  $J/\psi \rightarrow \eta X, X \rightarrow \gamma \eta'$  and  $J/\psi \rightarrow \eta' X, X \rightarrow \gamma \eta$  are constructed using the **covariant tensor formalism** 

• All kinematically allowed known resonances with  $0^{++}$ ,  $2^{++}$ ,  $4^{++}$  ( $\eta\eta'$ ) and  $1^{+-}$ ,  $1^{--}(\gamma\eta^{(\prime)})$  are considered

 $1^{-+}$  in  $\eta\eta'$  is also considered ( $\eta/\eta'$  not identical particle)

• An isoscalar resonance with exotic  $J^{PC} = 1^{-+}$ ,  $\eta_1(1855)$ , has been firstly observed with significance larger than  $19\sigma$ 

 $M = 1855 \pm 9^{+6}_{-1} \text{ MeV}/c^2$ ;  $\Gamma = (188 \pm 18^{+3}_{-8} \text{ MeV})$ 

 $\mathcal{B}(J/\psi \to \gamma \eta_1(1855) \to \gamma \eta \eta') = (2.70 \pm 0.41^{+0.16}_{-0.35}) \times 10^{-6}$ 

- Consistent with LQCD calculation for the 1<sup>-+</sup> hybrid (1.7~2.1 GeV/c<sup>2</sup>)
- > Hybrid? Molecule? Tetraquark? ... needs further study

PRL 129, 192002 (2022) PRL 130, 159901 (2023) PRD 106,072012 (2022) PRD 107,079901 (2023)



### **Observation of An Exotic Isoscalar State** $\eta_1(1855)$ (1<sup>-+</sup>) in $J/\psi \rightarrow \gamma \eta \eta'$

- Angular distribution expressed as an expansion in terms of Legendre polynomials **modelindependently**  $\langle Y_l^0 \rangle \equiv \sum_{i=1}^{N_k} W_i Y_l^0 (\cos \theta_{\eta}^i)$
- Related to the spin-0(S), spin-1(P), spin-2(D) amplitudes in  $\eta\eta'$  by :

$$\begin{split} \sqrt{4\pi} &< Y_0^0 >= S_0^2 + P_0^2 + P_1^2 + D_0^2 + D_1^2 + D_2^2 \\ \sqrt{4\pi} &< Y_1^0 >= 2S_0 P_0 cos\phi_{P_0} + \frac{2}{\sqrt{5}}(2P_0 D_0 \cos(\phi_{P_0} - \phi_{D_0}) + \sqrt{3}P_1 D_1 \cos(\phi_{P_1} - \phi_{D_1}) \\ \sqrt{4\pi} &< Y_2^0 >= \frac{1}{7\sqrt{5}}(14P_0^2 - 7P_1^2 + 10D_0^2 + 5D_1^2 - 10D_2^2) + 2S_0 D_0 cos\phi_{D_0} \\ \sqrt{4\pi} &< Y_3^0 >= \frac{6}{\sqrt{35}}(\sqrt{3}P_0 D_0 \cos(\phi_{P_0} - \phi_{D_0}) - P_1 D_1 \cos(\phi_{P_1} - \phi_{D_1})) \\ \sqrt{4\pi} &< Y_3^0 >= \frac{1}{7}(6D_0^2 - 4D_1^2 + D_2^2) \end{split}$$

- $< Y_1^0 >$  Indicates significant P-wave needed
- In ηη'system, only η<sub>1</sub> (1<sup>-+</sup>) contribute P-wave
   > Need for η<sub>1</sub>(1855)



# Discussions about $f_0(1500) \& f_0(1710)$

### • Decay properties :

- The pure glueball is a flavor singlet, the process  $G \rightarrow$ 
  - $\eta\eta'$  is significantly suppressed .

 $B(G \to \eta \eta')/B(G \to \pi \pi) < 0.04$ 

• Significant  $f_0(1500)$ 

 $\frac{B(f_0(1500) \to \eta \eta \prime)}{B(f_0(1500) \to \pi \pi)} = (1.66^{+0.42}_{-0.40}) \times 10^{-1}$ 

• Absence of  $f_0(1710)$ 

 $\frac{B(f_0(1710) \to \eta \eta')}{B(f_0(1710) \to \pi \pi)} < 2.87 \times 10^{-3} @90\% \text{ C. L.}$ 

Supports to the hypothesis that  $f_0(1710)$  overlaps with the ground state scalar  $(0^{++})$  glueball

Decay mode	Resonance	$M~({\rm MeV}/c^2)$	$\Gamma$ (MeV)	$M_{\rm PDG}~({\rm MeV}/c^2)$	$\Gamma_{\rm PDG}~({\rm MeV})$	B.F. $(\times 10^{-5})$	Sig.
	$f_0(1500)$	1506	112	1506	112	$1.81 \pm 0.11^{+0.19}_{-0.13}$	≫30σ
	$f_0(1810)$	1795	95	1795	95	$0.11{\pm}0.01^{+0.04}_{-0.03}$	11.1 <i>o</i>
	$f_0(2020)$	$2010{\pm}6^{+6}_{-4}$	$203{\pm}9^{+13}_{-11}$	1992	442	$2.28{\pm}0.12^{+0.29}_{-0.20}$	24.6 <i>σ</i>
$J/\psi \to \gamma X \to \gamma \eta \eta'$	$f_0(2330)$	$2312 \pm 7^{+7}_{-3}$	$65{\pm}10^{+3}_{-12}$	2314	144	$0.10{\pm}0.02^{+0.01}_{-0.02}$	13.2 <i>σ</i>
	$\eta_1(1855)$	$1855 \pm 9^{+6}_{-1}$	$188{\pm}18^{+3}_{-8}$	-	-	$0.27{\pm}0.04^{+0.02}_{-0.04}$	$21.4\sigma$
	$f_2(1565)$	1542	122	1542	122	$0.32{\pm}0.05^{+0.12}_{-0.02}$	$8.7\sigma$
	$f_2(2010)$	$2062{\pm}6^{+10}_{-7}$	$165{\pm}17^{+10}_{-5}$	2011	202	$0.71{\pm}0.06^{+0.10}_{-0.06}$	$13.4\sigma$
	$f_4(2050)$	2018	237	2018	237	$0.06{\pm}0.01^{+0.03}_{-0.01}$	$4.6\sigma$
	$0^{++}$ PHSP	-	-	-	-	$1.44{\pm}0.15^{+0.10}_{-0.20}$	15.7 <i>σ</i>
$J/\psi \to \eta' X \to \gamma \eta \eta'$	$h_1(1415)$	1416	90	1416	90	$0.08{\pm}0.01{}^{+0.01}_{-0.02}$	$10.2\sigma$
	$h_1(1595)$	1584	384	1584	384	$0.16{\pm}0.02^{+0.03}_{-0.01}$	9.9σ

## Partial Wave Analysis of $J/\psi \rightarrow \gamma \eta' \eta'$

#### PRD 105,072002 (2022)

### Observation of the three $f_0$ and one $f_2$

- **New state**  $f_0(2480)$
- Components of  $f_0(2020)$ 
  - flavor singlet

 $\frac{\Gamma(f_0(2020) \to \eta \eta \prime)}{\Gamma(f_0(2020) \to \eta \prime \eta \prime)} = 0.0148$ 

B.F.  $M(MeV/c^2)$ Γ(MeV) Significance  $(\sigma)$ Resonance  $436 \pm 4^{+46}_{-49}$  $f_0(2020)$  $(2.63 \pm 0.06^{+0.31}_{-0.46}) \times 10^{-4}$ ≫25  $1982 \pm 3^{+54}_{-0}$  $134 \pm 5^{+30}_{-9}$  $(6.09 \pm 0.64^{+4.00}_{-1.68}) \times 10^{-6}$  $f_0(2330)$  $2312 \pm 2^{+10}_{-0}$ 16.3  $^{-1.08}_{(8.18 \pm 1.77^{+3.73}_{-2.23}) \times 10^{-7}}$  $(4.69 \pm 0.80^{+0.74}_{-1.82}) \times 10^{-7}$  $(8.67 \pm 0.70^{+0.61}_{-1.67}) \times 10^{-6}$  $2470 \pm 4^{+4}_{-6}$  $75 \pm 9^{+11}_{-8}$ 5.2  $f_0(2480)$  $66 \pm 10^{+12}_{-10}$  $h_1(1415)$  $1384 \pm 6^{+9}_{-0}$ 5.3  $f_2(2340)$  $2346 \pm 8^{+22}_{-\epsilon}$  $332 \pm 14^{+26}_{-12}$ 16.10<sup>++</sup> PHSP  $(1.17 \pm 0.23^{+4.09}_{-0.70}) \times 10^{-5}$ 15.7



• New decay mode of tensor glueball candidate  $f_2(2340)$ 

$$\begin{split} & \mathsf{B}(\mathsf{J}/\psi\to\gamma f_2(2340)\to\gamma\eta\eta)=3.8^{+0.62}_{-0.65}{}^{+2.37}_{-2.07}\times10^{-5}_{\text{PRD 87,092009}} \\ & \mathsf{B}(\mathsf{J}/\psi\to\gamma f_2(2340)\to\gamma\phi\phi)=1.91\pm0.14^{+0.72}_{-0.73}\times10^{-4}_{-0.73}_{\text{PRD93,112011}} \\ & \mathsf{B}(\mathsf{J}/\psi\to\gamma f_2(2340)\to\gamma K_S K_S)=5.54^{+0.34+3.82}_{-0.40}\times10^{-5}_{-1.49}\times10^{-5}_{-0.40}_{-1.49}\times10^{-5}_{-0.65}_{-2.07}\times10^{-6}_{-0.65} \end{split}$$

### Mass Independent PWA(MI) explore the lineshape of $K_S^0 K_S^0 \pi^0$ invariant mass for the different decay modes, and minimize bias from model dynamics of the intermediate states

- Non-trivial  $0^{-+}$  line shape
- Mass Dependent PWA(MD) extract the resonance parameters of the intermediate states
  - Two resonances parameterization



Resonance

 $\eta(1405)$ 

 $\eta(1475)$ 

 $f_1(1285)$ 

 $f_1(1420)$ 

 $f_2(1525)$ 



### dominant spin-parity components





## A New State X(2600) Observed in $J/\psi \rightarrow \gamma \pi^+ \pi^- \eta'$

PRL 129, 042001 (2022)

- To study **X(2600)** parameters, a simultaneous fit to  $\eta' \pi^+ \pi^-$  and  $\pi^+ \pi^-$  is performed
- The structure in  $M(\pi^+\pi^-)$  well described with the interference between  $f_0(1500)$  and X(1540)

@ > 20σ	Mass $(MeV/c^2)$	Width $(MeV)$
$f_0(1500)$	$1492.5\pm3.6^{+2.4}_{-20.5}$	$107\pm9^{+21}_{-7}$
X(1540)	$1540.2\pm7.0^{+36.3}_{-6.1}$	$157 \pm 19^{+11}_{-77}$
X(2600)	$2618.3 \pm 2.0^{+16.3}_{-1.4}$	$195\pm5^{+26}_{-17}$
Case	$f_0(1500)$	X(1540)
Events	$24585 \pm 1689$	$21203 \pm 1456$
BF (×10 <sup>-</sup>	<sup>5</sup> ) $3.09 \pm 0.21^{+1.14}_{-0.77}$ 2	$2.69 \pm 0.19^{+0.38}_{-1.21}$

$$X(2600)$$
  
 $J^{PC} = 0^{-+} or 2^{-+}$ 



### **Observation of X(1835), X(2120) and X(2370) in** $J/\psi$ EM Dalitz Decays $J/\psi \rightarrow e^+ e^- \pi^+ \pi^- \eta'$ PRL 129, 022002

Events/(22.5 MeV/c<sup>2</sup>)





Measurement of the Transition Form Factor of  $J/\psi \rightarrow$  $e^+e^-X(1835)$ 

 $\succ$  Gives additional information of the internal structure of X(1835)

$$\frac{d\Gamma(J/\psi \to X(1835)e^+e^-)}{dq^2\Gamma(J/\psi \to X(1835)\gamma)} = |F(q^2)|^2 \times |QED(q^2)|$$
$$F(q^2) = \frac{1}{1 - q^2/\Lambda^2},$$
$$\Lambda = 1.75 \pm 0.29 \pm 0.05 \text{ GeV}/c^2$$

## Summary

...

- Exciting results from new  $J/\psi$  data are presented
  - Isoscalar 1<sup>-+</sup> exotic :  $\eta_1(1855)$
  - New state X(2600) in  $J/\psi$  radiative decays

- Data with unprecedented statistical accuracy from BESIII provides great opportunities to study QCD exotics. Will continue to run until  $\sim 2030$
- BESIII is in good status, inner detector upgrade in progress;
- ➢ BEPCII-U: 3x upgrade on luminosity

# Thank you for your attention!

Backup slide

### For comparison

Baseline set of amplitudes

PDG-optimized set of amplitudes



Narrow structure in  $\langle Y_1^0 \rangle$  cannot be described by resonances in  $\gamma \eta(\eta')$  $\gg \eta_1(1855) \rightarrow \eta \eta'$  needed

### Angular moments without contribution from D wave and F wave



The blue line becomes flat from a peak structure

### Further checks on the $1^{-+}$ state $\eta_1(1855)$

>  $\cos\theta_{\eta}$  distributions in different M( $\eta\eta'$ ) regions



✓ Clear asymmetry in the region [1.7,2.0]GeV/ $c^2$ , largely due to  $\eta_1(1855)$  signal

### Scan of additional resonance (based on **Baseline set of amplitudes**)



 $\square$  If two 1<sup>-+</sup> are introduced in PWA

- $\eta_1(1855)$ : M=1873 ± 7 MeV/ $c^2$ ,  $\Gamma = 225 \pm 7$  MeV, Br is  $(2.6 \pm 0.2) \times 10^{-6}$ , significance is 15.7 $\sigma$
- $\eta_1(\text{additional})$ : M=2152 ± 3 MeV/ $c^2$ ,  $\Gamma = 134 \pm 2$  MeV, Br is  $(3 \pm 1) \times 10^{-7}$ , significance is 4.6 $\sigma$

> The  $\eta_1$  (additional) is not included in PWA, because of small fraction (0.4%), and significance < 5 $\sigma$  22

## PWA of $J/\psi \rightarrow \gamma \eta \eta$

Events / 0.020 GeV/c<sup>2</sup>

(a)

1.5

 $2.25 \times 10^8 J/\psi$  at BESIII

### PRD 87, 092009 (2013)



lesonance	${\rm Mass}({\rm MeV}/c^2)$	${\rm Width}({\rm MeV}/c^2)$	$\mathcal{B}(J/\psi \to \gamma X \to \gamma \eta \eta)$	Significance
$f_0(1500)$	$1468^{+14+23}_{-15-74}$	$136^{+41+28}_{-26-100}$	$(1.65^{+0.26+0.51}_{-0.31-1.40}) \times 10^{-5}$	$8.2 \sigma$
$f_0(1710)$	$1759 \pm 6^{+14}_{-25}$	$172 \pm 10^{+32}_{-16}$	$(2.35^{+0.13+1.24}_{-0.11-0.74}) \times 10^{-4}$	25.0 $\sigma$
$f_0(2100)$	$2081{\pm}13^{+24}_{-36}$	$273^{+27+70}_{-24-23}$	$(1.13^{+0.09+0.64}_{-0.10-0.28}) \times 10^{-4}$	13.9 $\sigma$
$f_{2}^{'}(1525)$	$1513 \pm 5^{+4}_{-10}$	$75_{-10-8}^{+12+16}$	$(3.42^{+0.43+1.37}_{-0.51-1.30}) \times 10^{-5}$	11.0 $\sigma$
$f_2(1810)$	$1822^{+29+66}_{-24-57}$	$229_{-42-155}^{+52+88}$	$(5.40^{+0.60+3.42}_{-0.67-2.35}) \times 10^{-5}$	6.4 $\sigma$
$f_2(2340)$	$2362_{-30-63}^{+31+140}$	$334_{-54-100}^{+62+165}$	$(5.60^{+0.62+2.37}_{-0.65-2.07}) \times 10^{-5}$	7.6 $\sigma$

PWA of  $J/\psi \rightarrow \gamma K_S^0 K_S^0$ 

1311M  $J/\psi$  at BESIII

#### PRD 98, 072003 (2018)



χ<sup>2</sup>/nbin=1.32

- Global Fit

-4

-3

-2

+ Data

-0.2

cos(θ<sub>Ks</sub>)

0.2 0.4 0.6 0.8

8000 7000

6000

2000

1000

0

-2 -4

-1

-0.8 -0.6 -0.4

Pull



0

¢κs

+++

3

Resonance	$M (\text{MeV}/c^2)$	$M_{\rm PDG}~({\rm MeV}/c^2)$	$\Gamma (\text{MeV}/c^2)$	$\Gamma_{\rm PDG} \ ({\rm MeV}/c^2)$	Branching fraction	Significance
$K^{*}(892)$	896	$895.81 {\pm} 0.19$	48	$47.4 {\pm} 0.6$	$(6.28^{+0.16}_{-0.17}, 0.59}_{-0.17}) \times 10^{-6}$	$35\sigma$
$K_1(1270)$	1272	$1272\pm7$	90	$90{\pm}20$	$(8.54^{+1.07}_{-1.20}) \times 10^{-7}$	$16\sigma$
$f_0(1370)$	$1350 \pm 9^{+12}_{-2}$	1200 to 1500	$231 \pm 21^{+28}_{-48}$	200 to 500	$(1.07^{+0.08}_{-0.07}, 0.34}^{+0.08}) \times 10^{-5}$	$25\sigma$
$f_0(1500)$	1505	$1504 \pm 6$	109	$109 \pm 7$	$(1.59^{+0.16}_{-0.16}^{+0.18}) \times 10^{-5}$	$23\sigma$
$f_0(1710)$	$1765 \pm 2^{+1}_{-1}$	$1723^{+6}_{-5}$	$146 \pm 3^{+7}_{-1}$	$139 \pm 8$	$(2.00^{+0.03}_{-0.02}{}^{+0.31}_{-0.10})\times 10^{-4}$	$\gg 35\sigma$
$f_0(1790)$	$1870\pm7^{+2}_{-3}$	-	$146 \pm 14^{+7}_{-15}$	-	$(1.11^{+0.06}_{-0.06}^{+0.19}_{-0.32}) \times 10^{-5}$	$24\sigma$
$f_0(2200)$	$2184\pm5^{+4}_{-2}$	$2189 \pm 13$	$364 \pm 9^{+4}_{-7}$	$238 \pm 50$	$(2.72^{+0.08}_{-0.06}^{+0.17})\times 10^{-4}$	$\gg 35\sigma$
$f_0(2330)$	$2411\pm10\pm7$	-	$349 \pm 18^{+23}_{-1}$	-	$(4.95^{+0.21}_{-0.21}{}^{+0.66}_{-0.72})\times 10^{-5}$	$35\sigma$
$f_2(1270)$	1275	$1275.5 \pm 0.8$	185	$186.7^{+2.2}_{-2.5}$	$(2.58^{+0.08}_{-0.09}, 0.20}_{-0.09}) \times 10^{-5}$	$33\sigma$
$f'_2(1525)$	$1516 \pm 1$	$1525\pm5$	$75 \pm 1 \pm 1$	$73^{+6}_{-5}$	$(7.99^{+0.03}_{-0.04}^{+0.69}) \times 10^{-5}$	$\gg 35\sigma$
$f_2(2340)$	$2233 \pm 34^{+9}_{-25}$	$2345^{+50}_{-40}$	$507 \pm 37^{+18}_{-21}$	$322^{+70}_{-60}$	$(5.54^{+0.34}_{-0.40}) \times 10^{-5}$	$26\sigma$
0 <sup>++</sup> PHSP	-	-	-	-	$(1.85^{+0.05}_{-0.05}{}^{+0.68}_{-0.26})\times 10^{-5}$	$26\sigma$
$2^{++}$ PHSP	-	-	-	-	$(5.73^{+0.99+4.18}_{-1.00-3.74}) \times 10^{-5}$	$13\sigma$



FIG. 1. Two Regge trajectories for the isovector and isoscalar scalar states. Here, the red triangles denote the  $a_0(1817)$  and X(1812), while the solid points and empty circles are the experimental and predicted states, respectively. Except the  $a_0(1817)$ , X(1812) and predicted  $a_0(2115)$ ,  $f_0(1450)$ , the masses of the other scalar states are taken from PDG. Here, the established states are marked by the black solid points, while the predicted states are denoted by the black circles. The error bars present total experimental uncertainties.

# Detail plots of $J/\psi \rightarrow \gamma K_s^0 K_s^0 \pi^0$



Mass distributions and Dalitz plots





dominant decay modes for the axial vector component

# Detail plots of $J/\psi \rightarrow \gamma K_s^0 K_s^0 \pi^0$



Resonance	Decay Mode	Event selection	Breit-Wigner formula	Resonance parameters	Extra components	Total
m(1405)	$J/\psi \to \gamma \eta (1405) \to \gamma K^0_S (K^0_S \pi^0)_{\text{P-wave}} \to \gamma K^0_S K^0_S \pi^0$	$\pm 4.0$	$^{+28.1}_{-56.8}$	+20.0	$^{+2.8}_{-7.8}$	$^{+34.8}_{-57.5}$
$\eta(1405) =$	$J/\psi \to \gamma \eta (1405) \to \gamma (K^0_S K^0_S)_{\text{S-wave}} \pi^0 \to \gamma K^0_S K^0_S \pi^0$	$\pm 4.0$	$^{+19.0}_{-7.5}$	+53.1	$^{+6.1}_{-10.0}$	$^{+56.9}_{-13.1}$
n(1475)	$J/\psi \to \gamma \eta (1475) \to \gamma K^0_S (K^0_S \pi^0)_{\text{P-wave}} \to \gamma K^0_S K^0_S \pi^0$	$\pm 4.0$	$^{+59.4}_{-42.5}$	+8.9	$^{+5.8}_{-3.1}$	$^{+60.5}_{-42.8}$
$\eta(1475) =$	$J/\psi \to \gamma \eta (1475) \to \gamma (K^0_S K^0_S)_{\text{S-wave}} \pi^0 \to \gamma K^0_S K^0_S \pi^0$	$\pm 4.0$	$^{+5.5}_{-7.4}$	-13.7	$^{+7.8}_{-4.3}$	$^{+10.3}_{-16.6}$
$f_1(1285)$	$J/\psi \rightarrow \gamma f_1(1285) \rightarrow \gamma a_0(980)^0 \pi^0 \rightarrow \gamma K_S^0 K_S^0 \pi^0$	$\pm 4.0$	$^{+1.0}_{-11.2}$	+4.3	$^{+39.6}_{-2.5}$	$^{+40.0}_{-12.2}$
£.(1420)	$J/\psi \rightarrow \gamma f_1(1420) \rightarrow \gamma K^*(892)^0 K^0_S \rightarrow \gamma K^0_S K^0_S \pi^0$	$\pm 4.0$	$^{+2.5}_{-16.8}$	+6.4	$^{+6.1}_{-1.2}$	$^{+10.0}_{-17.3}$
J1(1420)	$J/\psi \rightarrow \gamma f_1(1420) \rightarrow \gamma a_0(980)^0 \pi^0 \rightarrow \gamma K_S^0 K_S^0 \pi^0$	$\pm 4.0$	$^{+15.8}_{-41.8}$	+7.8	+47.8	$^{+51.1}_{-42.0}$
$f_2(1525)$	$J/\psi \rightarrow \gamma f_2(1525) \rightarrow \gamma K^*(892)^0 K^0_S \rightarrow \gamma K^0_S K^0_S \pi^0$	$\pm 4.0$	$^{+15.3}_{-4.7}$	-3.3	$^{+1.9}_{-0.5}$	$^{+15.9}_{-7.0}$

### MD PWA results

X(18xx) between 1.8-1.9GeV



## Other results on X(1835)

### PRL 117, 042002 (2016)



 $1.09 \times 10^9 J/\psi$  at BESIII  $J/\psi \rightarrow \gamma \pi^+ \pi^- \eta'$ 

### Significant distortion of the line $\pi^+\pi^-\eta'$ shape near the ppbar mass threshold

Two fit models are taken into account and both support the existence of a  $p\bar{p}$  moleculelike or bound state

### PRD 97,051101(R) (2018)



 $1.3 \times 10^9 J/\psi$  at BESIII

J/ψ → γγφ: two structures corresponding to η(1475) and X(1835) are observed

• X(1835) and  $\eta(1475)$ :  $J^{PC} = 0^{-+}$  assignment favored

- Sizalbe s $\bar{s}$  component in X(1835)
  - More complicated than a pure  $N\overline{N}$  state

Solution Resonance	$m_R ({ m MeV}/c^2)$	$\Gamma ({ m MeV})$	$B(10^{-6})$
$\eta(1475)$	$1477\pm7\pm13$	$118\pm22\pm17$	$7.03 \pm 0.92 \pm 0.91$
(Destr. Int.) $X(1835)$	$1839\pm26\pm26$	$175\pm57\pm25$	$1.77 \pm 0.35 \pm 0.25$
$\eta(1475)$	$1477\pm7\pm13$	$118\pm22\pm17$	$10.36 \pm 1.51 \pm 1.54$
(Constr. Int.) $X(1835)$	$1839\pm26\pm26$	$175\pm57\pm25$	$8.09 \pm 1.99 \pm 1.36$

## Search for X(1835) in other decay modes



 $1.3 \times 10^9 J/\psi$  at BESIII  $J/\psi \rightarrow \omega \eta' \pi^+ \pi^-$ 

No obvious sign of X(1835)'s existence

$$\begin{split} B(J/\psi \to \omega \eta' \pi^+ \pi^-) &= (1.12 \pm 0.02 \pm 0.13) \times 10^{-3} \\ B(J/\psi \to \omega X(1835), X(1835) \to \eta' \pi^+ \pi^-) < 6.2 \times 10^{-5} \end{split}$$

@90%C.L.

# First Observation of X(2370) in $J/\psi \rightarrow \gamma K \overline{K} \eta'$

 $1.3 \times 10^9 J/\psi$  at BESIII

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- X(2120) and X(2370) states observed in the  $\eta' \pi^+ \pi^-$  invariant mass spectra(PRL106,072002)
- The X(2370) measured mass is consistent with the pseudoscalar glueball candidate predicted by LQCD calculation(PRD73,014516)



No evidence of X(2120)  $B(J/\psi \rightarrow \gamma X(2120) \rightarrow \gamma K^+ K^- \eta') < 1.49 \times 10^{-5}$  $B(J/\psi \rightarrow \gamma X(2120) \rightarrow \gamma K_S^0 K_S^0 \eta') < 6.38 \times 10^{-5}$ 

Clear X(2370) signal observed with significance of 8.3 $\sigma$   $M_{X(2370)} = 2341.6 \pm 6.5 \pm 5.7 \text{ MeV/c}^2$ ;  $\Gamma_{X(2370)} = 117 \pm 10 \pm 8 \text{ MeV}$   $B(J/\psi \rightarrow \gamma X(2370) \rightarrow \gamma K^+ K^- \eta') < 1.49 \times 10^{-5}$  $B(J/\psi \rightarrow \gamma X(2370) \rightarrow \gamma K_S^0 K_S^0 \eta') < 6.38 \times 10^{-5}$ 

## First Observation of X(2370) in $J/\psi \rightarrow \gamma K K \eta'$

200

150

EPJC 80, 746 (2020)  $1.3 \times 10^9 J/\psi$  at BESIII  $J/\psi \rightarrow \gamma K^{\dagger} K \eta', \eta' \rightarrow \pi^{\dagger} \pi \eta,$ η→γγ (a) 500  $\rightarrow \pi^{+}\pi^{-}$ K'K n', n' Events/(0.01GeV/c<sup>2</sup>) Chebychev Chebychev Data Data ---- PHSP PHSP 400 Fit result Fit result ····· Total bkg ····· Total bkg 🗱 Signal X(2370) Signal X(2370) 300  $J/\psi \rightarrow K^{+}K^{-}\eta' + c.c$  $J/\psi \rightarrow K^{*+} K^{-} \eta' + c.c.$ n' sideband n' sideband 200 100

Events/(0.01GeV/c<sup>2</sup>) 100 50 02 2 2.2 2.3 2.4 2.5 2.6 2.7 2.5 2.6 2.2 2.3 2.4 2.7 2.1 2.1  $M_{K^{+}K^{-}n^{\prime}}$  (GeV/c<sup>2</sup>)  $M_{K^{+}K^{-}\eta^{+}}$  (GeV/c<sup>2</sup>) 30 90 (d) (c)  $J/\psi \rightarrow \gamma K_{s}^{o}K_{s}^{o}\eta', \eta' \rightarrow \pi^{*}\pi^{*}\eta, \eta \rightarrow \gamma \gamma$  $\rightarrow \pi^{+}\pi^{-}$  $\rightarrow \gamma K_{c}^{\circ} K_{c}^{\circ} \eta', \eta' \rightarrow$ 80 Events/(0.01GeV/c<sup>2</sup>) Events/(0.01GeV/c<sup>2</sup>) 25 Chebychev Chebychev Data Data ---- PHSP PHSP Fit result Fit result ······ Total bkg ······ Total bkg 20 60 🔆 Signal X(2370) Signal X(2370) 50 η' sideband n' sideband 15 10 02 <sup>0</sup>2 2.2 2.5 2.6 2.7 2.3 2.5 2.6 2.7 2.1 2.3 2.4 2.2 2.4 2.1  $M_{K_s^0 K_s^0 \eta'}$  (GeV/c<sup>2</sup>)  $M_{K_s^0 K_s^0 \eta'}$  (GeV/c<sup>2</sup>)

(b)

# Search for X(2370) in $J/\psi \rightarrow \gamma \eta \eta \eta'$

 $1.3 \times 10^9 J/\psi$  at BESIII

### PRD 103, 012009 (2021)



No evidence of X(2120)  $B(J/\psi \rightarrow \gamma X(2370) \rightarrow \gamma \eta \eta \eta') < 9.2 \times 10^{-6}$ First Observation in the  $\eta \eta \eta'$  invariant mass spectra  $B(J/\psi \rightarrow \gamma \eta_c) \cdot B(\eta_c \rightarrow \eta \eta \eta') = (4.86 \pm 0.62 \pm 0.45) \times 10^{-5}$