## $11^{\text {th }}$ International Workshop on Charm Physics（CHARM 2023）



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Jul．17－21，2023，Siegen，Germany

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## Physics Motivation



The decays of $\mathrm{J} / \psi$ and $\psi(3686)$ to hadronic final states are used to test the " $12 \%$ rule".

$$
Q_{h}=\frac{\mathcal{B}_{\psi(3686) \rightarrow h}}{\mathcal{B}_{J / \psi \rightarrow h}} \approx \frac{\mathcal{B}_{\psi(3686) \rightarrow e^{+} e^{-}}}{\mathcal{B}_{J / \psi \rightarrow e^{+} e^{-}}} \approx 13.3 \%
$$

The baryon/meson pair decays of $\chi_{\mathbf{c} 0,1,2}$ are essential to test pQCD, e.g. helicity selection rule (HSR) . In addition, it is used to test color-octet mechanism (COM).

* The search of $\psi(3770)$ non-DD decay modes is helpful to understand not only the properties of $\psi(3770)$ and $\psi(3686)$, e.g., 2S-1D mixing, but also for exotic charmoniumlike (XYZ) states.

$$
\text { Kind reminder: } \quad \psi(3686) \equiv \psi(2 S) \equiv \psi^{\prime}
$$

## Charmonium Spectrum



# Traditional charmonium 

states are named as
$\eta_{c}, \psi, h_{c}, \chi_{c} \ldots \ldots$
charmoniumlike states are named with "XYZ" states,


## BESIII Detector



## BESIII Data Collections


$\checkmark$ BESIII detector has collected the largest data samples in the $\psi$ energy region in the world.
$\checkmark \sim 30 \%$ of $\psi(3686)$ produces $\chi_{c J}$ in its EM radiative transition decay.

## Recent Physics Highlights on $\chi_{c J}$ and $\psi$ decays

## BESIII

- Measurement of $\mathrm{B}\left[\psi(3686) \rightarrow \phi K_{S}^{0} K_{S}^{0}\right]$
- Helicity amplitude analysis of $\chi_{c J} \rightarrow \phi \phi$
- Observation of $\chi_{c J} \rightarrow \Omega^{-} \bar{\Omega}^{+}$
- Observation of $\psi(3770) \rightarrow \eta J / \psi$


## Measurement of $\mathrm{B}\left[\psi(3686) \rightarrow \phi K_{S}^{0} K_{S}^{0}\right]$






- The QED background is investigated using the continuum data @ 3.65 GeV by luminosity normalization. It is determined to be $\bar{N}_{Q E D}=108 \pm 5$
- The interference between $\psi(3686)$ and continuum production is considered by fitting to the cross section around $\psi(3686)$. The number of events due to interference is determined to be $N_{\text {inter }}=228 \pm 24$
$\mathcal{B}_{\psi(3686) \rightarrow \phi K_{S}^{0} K_{S}^{0}}=\frac{N_{\text {net }}^{\psi(3686)}-\bar{N}_{\text {QED }}-N_{\text {inter }}}{N_{\psi(3686)} \cdot \varepsilon \cdot \mathcal{B}_{\phi \rightarrow K^{+} K^{-}} \cdot \mathcal{B}_{K_{S}^{0} \rightarrow \pi^{+} \pi^{-}}^{2}}=(3.53 \pm 0.20 \pm 0.21) \times 10^{-5}$
$Q_{\psi(3686) \rightarrow \phi K_{S}^{0} K_{S}^{0}}=(6.0 \pm 1.6) \%$, which is suppressed relative to " $12 \%$ rule"


## Helicity amplitude analysis of $\chi_{c J} \rightarrow \phi \phi(1)$

> The asymptotic behavior of the BF of charmonium state $\psi(\lambda)$ decaying into light hadron $h_{1}\left(\lambda_{1}\right)$ and $h_{2}\left(\lambda_{2}\right)$ is evaluated in pQCD calculation


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(a)
where, $\lambda$ denotes the helicity of hadrons, $\Lambda_{\mathrm{QCD}}$ denotes the QCD scale parameter
$>$ The $\chi_{\mathrm{c} 1}$ decay rate is suppressed relative to $\chi_{\mathrm{c} 0,2}$ due to helicity selection rule (HSR) according to pQCD prediction, which is not consistent with the experimental measurement.


Definitions of helicity angle

## Helicity amplitude analysis of $\chi_{c J} \rightarrow \phi \phi$ (2)

| Decay Mode | Helicity Angle | Amplitude |
| :---: | :---: | :---: |
| $\psi(3686)(M) \rightarrow R_{i}\left(\lambda_{R}\right) \gamma\left(\lambda_{\gamma}\right)$ | $\theta_{0}$ | $A_{\lambda_{\gamma}, \lambda_{R}}^{1} D_{M, \lambda_{R}-\lambda_{\gamma}}^{1 *}\left(0, \theta_{0}, 0\right)$ |
| $R_{i}\left(\lambda_{R}\right) \rightarrow \phi\left(\lambda_{1}\right) \phi\left(\lambda_{2}\right)$ | $\theta_{1}, \phi_{1}$ | $F_{\lambda_{1}, \lambda_{2}}^{J} D_{\lambda_{R}, \lambda_{1}-\lambda_{2}}^{J_{2}^{*}}\left(\phi_{1}, \theta_{0}, 0\right)$ |
| $\phi\left(\lambda_{1}\right) \rightarrow K^{+}\left(0^{-}\right) K^{-}\left(0^{-}\right)$ | $\theta_{2}, \phi_{2}$ | $B_{0,0}^{1} D_{\lambda_{1}, 0}^{1 *}\left(\phi_{2}, \theta_{2}, 0\right)$ |
| $\phi\left(\lambda_{2}\right) \rightarrow K^{+}\left(0^{-}\right) K^{-}\left(0^{-}\right)$ | $\theta_{3}, \phi_{3}$ | $B_{0,0}^{1} D_{\lambda_{2}, 0}^{1 *}\left(\phi_{3}, \theta_{3}, 0\right)$ |

Table 2. Definitions of helicity angles and amplitudes of sequential decays.




For $\chi_{c 1} \rightarrow \phi \phi$, the ratios of amplitude moduli (T/L) are determined to be

$$
u_{1}=\left|F_{1,0}^{1}\right| F_{0,1}^{1} \mid=1.05 \pm 0.05 \text { and } u_{2}=\left|F_{1,1}^{1} / F_{1,0}^{1}\right|=0.07 \pm 0.04 .
$$

The results are consistent with the expectation of identical particle symmetry and parity conservation.

## Helicity amplitude analysis of $\chi_{c J} \rightarrow \phi \phi$ (3)

## For $\chi_{c 0}$

$$
\begin{array}{||l|l|l|}
\hline\left|F_{1,1}^{0}\right| /\left|F_{0,0}^{0}\right|=0.299 \pm 0.003 \pm 0.019 \\
\hline
\end{array}
$$

$$
\left|F_{0,1}^{2}\right| /\left|F_{0,0}^{2}\right|=1.265 \pm 0.054 \pm 0.079
$$

For $\chi_{\mathbf{c} 2} \longrightarrow\left|F_{1,-1}^{2}\right| /\left|F_{0,0}^{2}\right|=1.450 \pm 0.097 \pm 0.104$

$$
\left|F_{1,1}^{2}\right| /\left|F_{0,0}^{2}\right|=0.808 \pm 0.051 \pm 0.009
$$

| Decay Mode | 2011 BESIII $[8]$ | this work | PDG value $[22]$ |
| :---: | :---: | :---: | :---: |
| $\mathcal{B}\left[\chi_{c 0} \rightarrow \phi \phi\right]\left(\times 10^{-4}\right)$ | $7.8 \pm 0.4 \pm 0.8$ | $8.59 \pm 0.27 \pm 0.20$ | $8.0 \pm 0.7$ |
| $\mathcal{B}\left[\chi_{c 1} \rightarrow \phi \phi\right]\left(\times 10^{-4}\right)$ | $4.1 \pm 0.3 \pm 0.5$ | $4.26 \pm 0.13 \pm 0.15$ | $4.2 \pm 0.5$ |
| $\mathcal{B}\left[\chi_{c 1} \rightarrow \phi \phi\right]\left(\times 10^{-4}\right)$ | $10.7 \pm 0.4 \pm 1.2$ | $12.67 \pm 0.28 \pm 0.33$ | $10.6 \pm 0.9$ |

Table 6. Comparsion of measured branching fractions $(\mathcal{B})$.
This measurements are the most precise!

## Observation of $\chi_{c J} \rightarrow \Omega^{-} \bar{\Omega}^{+}$

$>$ This decay is unique due to the presence of three pairs of strange quarks in the final state.
$>\Omega$ is the only baryon of the ground-state decuplet decaying through weak interaction. Therefore, it has a long lifetime.
> Single baryon reconstruction technique is applied to enhance the detection efficiency. i.e., we reconstruct $\Omega^{-}$, and $\bar{\Omega}^{+}$is tagged by the recoil mass.

| Mode | $N_{\chi_{c J}}^{\text {obs }}$ | $\epsilon_{\chi_{c J}}(\%)$ | Sig. $(\sigma)$ | $\mathcal{B}\left(\times 10^{-5}\right)$ |
| :--- | ---: | :---: | :---: | :---: |
| $\chi_{c 0}$ | $284 \pm 44$ | 3.05 | 5.6 | $3.51 \pm 0.54$ |
| $\chi_{c 1}$ | $277 \pm 42$ | 7.02 | 6.4 | $1.49 \pm 0.23$ |
| $\chi_{c 2}$ | $1038 \pm 56$ | 8.91 | 18 | $4.52 \pm 0.24$ |

> The measured $\chi_{\mathrm{c} 0}$ decay branching fraction is obviously small than those decaying into with spin $1 / 2$ and $3 / 2$.

PRD 107, 092004 (2023)



## Observation of $\psi(3770) \rightarrow \eta J / \psi$




> The Born cross section at 3.773 GeV is determined by

$$
\begin{aligned}
& \sigma^{B}\left(e^{+} e^{-} \rightarrow \eta J / \psi\right)=\frac{N^{\text {obs }}}{\mathcal{L} \cdot\left(1+\delta^{\mathrm{ISR}}\right) \cdot\left(1+\delta^{\mathrm{VP}}\right) \cdot \varepsilon \cdot \mathcal{B} r}, \\
&=\left(8.88 \pm 0.87_{\text {stat }} \pm 0.42_{\text {sys. }}\right) \mathrm{pb} \\
& \mathcal{B}(\psi(3770) \rightarrow \eta \mathrm{J} / \Psi)=\left(8.7 \pm 1.0_{\text {stat }} \pm 0.8_{\text {sys. }}\right) \times 10^{-4}
\end{aligned}
$$ is determined by the global fit to line-shape without considering the interference with the other processes, which is close to CLEO's result.


$>$ If consider the interference, the BF varies between $\left.11.2 \pm 5 . \mathbf{8}_{\text {stat }} \pm 1.1_{\text {sys. }}\right) \times \mathbf{1 0}^{-4}$ to $\left.11.6 \pm 6.0_{\text {stat }} \pm 1.1_{\text {sys. }}\right) \times 10^{-4}$ PRD 107, L091101 (2023)
$>$ The statistical significance is above $7 \sigma$ whether the interference is considered or not

## Summary

* The branching fraction of $\psi(3686) \rightarrow \phi K_{S}^{0} K_{S}^{0}$ is observed for the first time, and the ratio $\mathcal{B}\left(\psi(3686) \rightarrow \phi K_{S}^{0} K_{S}^{0}\right) / \mathcal{B}\left(J / \psi \rightarrow \phi K_{S}^{0} K_{S}^{0}\right)$ is found to be suppressed relative to " $12 \%$ rule".
$\star$ The branching fractions for $\chi_{c J} \rightarrow \phi \phi$ are significantly improved, and the corresponding polarization parameters are measured for the first time.
$\chi_{c J} \rightarrow \Omega^{-} \bar{\Omega}^{+}$are firstly observed and the corresponding branching fractions are measured.
* A new non-DD decay mode, $\psi(3770) \rightarrow \eta \mathrm{J} / \psi$, is confirmed, and the corresponding branching fraction is measured with and without considering the interference.
Much more results will be presented in the future.

> Thank you!

## Backup

## $\psi(3686) \rightarrow \phi K_{S}^{0} K_{S}^{0}$



FIG. 3. Dalitz plots of $M_{\phi K_{S 1}^{0}}^{2}$ vs. $M_{\phi K_{S 2}^{0}}^{2}$ for the accepted candidates in (a) data and (b) BODY3 signal MC events.

$$
\sigma^{\mathrm{dressed}}(s)=\left|A_{\mathrm{cont}}(s)+A_{\mathrm{res}}(s) \times e^{i \varphi}\right|^{2}
$$


(a)

(b)

## Helicity amplitude analysis of $\chi_{c J} \rightarrow \phi \phi$

Then the joint amplitude for the sequential process is obtained by

$$
\begin{aligned}
\mathcal{M}\left(R_{i}\right)= & \frac{1}{2} \sum_{M, \lambda_{R}, \lambda_{1}, \lambda_{2}} A_{\lambda_{R}, \lambda_{\gamma}}^{1} D_{M, \lambda_{R}-\lambda_{\gamma}}^{1 *}\left(0, \theta_{0}, 0\right) F_{\lambda_{1}, \lambda_{2}}^{J} D_{\lambda_{R}, \lambda_{1}-\lambda_{2}}^{J *}\left(\phi_{1}, \theta_{0}, 0\right) \\
& \times B_{0,0}^{1} D_{\lambda_{1}, 0}^{1 *}\left(\phi_{2}, \theta_{2}, 0\right) B_{0,0}^{1} D_{\lambda_{2}, 0}^{1 *}\left(\phi_{3}, \theta_{3}, 0\right) B W\left(m_{\phi \phi}, m_{i}, \Gamma_{i}\right),
\end{aligned}
$$

## $\psi(3770) \rightarrow \eta J / \psi$

$$
\begin{aligned}
\sigma_{\mathrm{co}}= & \mid C \cdot \sqrt{\Phi(s)}+e^{i \phi_{1}} \mathrm{BW}_{\psi(3770)}+e^{i \phi_{2}} \mathrm{BW}_{\psi(4040)} \\
& +e^{i \phi_{3}} \mathrm{BW}_{Y(4230)}+\left.e^{i \phi_{4}} \mathrm{BW}_{Y(4390)}\right|^{2} \\
\sigma_{\text {inco }}= & \left|\mathrm{BW}_{\psi(3770)}\right|^{2}+\mid C \cdot \sqrt{\Phi(s)}+e^{i \phi_{2}} \mathrm{BW}_{\psi(4040)} \\
& +e^{i \phi_{3}} \mathrm{BW}_{Y(4230)}+\left.e^{i \phi_{4}} \mathrm{BW}_{Y(4390)}\right|^{2}
\end{aligned}
$$

TABLE III. Fitting results of the $e^{+} e^{-} \rightarrow \eta J / \psi$ decay. The uncertainties on the branching fractions and $\phi$ are statistical and systematic. The $C_{0}$ of the four solutions in the coherent fit are the same.

|  | Coherent fit |  |  |  |  |
| :--- | :---: | :---: | :---: | :---: | :---: |
| Parameters | Solution1 | Solution2 | Solution3 | Solution4 | Incoherent fit |
| $M_{1}\left(\mathrm{MeV} / c^{2}\right)$ |  | 3773.7 (fixed) | 3773.7 (fixed) |  |  |
| $\Gamma_{1}(\mathrm{MeV})$ |  | 27.2 (fixed) |  | 27.2 (fixed) |  |
| $C_{0}$ |  | $13.3 \pm 1.9$ | $11.0 \pm 1.6$ |  |  |
| $\mathcal{B} r_{1}\left(\times 10^{-4}\right)$ | $11.3 \pm 5.9 \pm 1.1$ | $11.6 \pm 6.0 \pm 1.1$ | $11.2 \pm 5.8 \pm 1.1$ | $11.5 \pm 6.0 \pm 1.1$ | $8.7 \pm 1.0 \pm 0.8$ |
| $\phi_{1}(\mathrm{rad})$ | $3.9 \pm 0.6 \pm 0.07$ | $4.2 \pm 0.6 \pm 0.09$ | $3.7 \pm 0.6 \pm 0.05$ | $4.1 \pm 0.6 \pm 0.08$ |  |

## $\psi(3770) \rightarrow \eta J / \psi$

TABLE IV. Relative systematic uncertainties in percent on the branching fraction of $\psi(3770) \rightarrow \eta J / \psi$.

| Source | Coherent fit |  |  |  |  |  |  |  | Incoherent fit |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
|  | Solution1 |  | Solution2 |  | Solution3 |  | Solution4 |  |  |
|  | Br | $\phi$ | Br | $\phi$ | Br | $\phi$ | Br | $\phi$ | Br |
| Center-of-mass energy | 3.3 | 0.2 | 3.5 | 0.1 | 3.4 | 0.1 | 3.2 | 0.2 | 2.1 |
| Energy spread | 0.8 | 0.1 | 0.9 | 0.1 | 0.9 | 0.1 | 0.8 | 0.1 | 1.0 |
| $\psi(3770)$ mass | 1.5 | 0.1 | 1.4 | 0.1 | 1.5 | 0.1 | 1.4 | 0.1 | 0.9 |
| $\psi(3770)$ width | 4.2 | 0.1 | 4.2 | 0.1 | 4.1 | 0.1 | 4.1 | 0.1 | 3.6 |
| $\psi(3770) \Gamma_{e^{+} e^{-}}$ | 6.9 |  | 6.9 |  | 6.9 |  | 6.9 |  | 6.9 |
| $\psi(4040)$ mass | 0.5 | 0.4 | 0.2 | 0.1 | 0.4 | 0.3 | 0.3 | 0.1 | 0.4 |
| $\psi(4040)$ width | 0.7 | 0.3 | 1.0 | 0.8 | 1.1 | 0.5 | 0.8 | 0.8 | 0.7 |
| Continuum term | 0.9 | 1.8 | 0.9 | 2.0 | 0.9 | 1.3 | 1.0 | 1.6 | 1.3 |
| Correlated systematic uncertainties | 3.0 |  | 3.0 |  | 3.0 |  | 3.0 |  | 3.0 |
| Uncorrelated systematic uncertainties | 0.2 | 0.1 | 0.1 | 0.1 | 0.1 | 0.1 | 2.0 | 0.5 | 0.1 |
| Total | 9.5 | 1.9 | 9.5 | 2.2 | 9.5 | 1.4 | 9.6 | 1.9 | 8.8 |

