### CHARMED MESON AND BARYON SPECTROSCOPY Selected recent experimental results

#### Patrick Spradlin on behalf of the LHCb collaboration with input from representatives of Belle and BESIII





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**OPEN CHARM SPECTROSCOPY** 

### OUTLINE



## New baryon resonances

NEW CHARMED BARYON AT BELLE: PRL 130 (2023) 3, 031901



ANALYSIS OF  $\Sigma_{c}^{++(0)}\pi^{-(+)}$  IN  $\overline{B}^{0}$  DECAYS Belle: Phys.Rev.Lett. 130 (2023) 3, 031901 (arXiv:2206.08822 [hep-ex])

Part of an analysis of  $\overline{B}{}^0 \rightarrow \Lambda_c^+ \pi^+ \pi^- \overline{\rho}$ 

- Full Belle  $\Upsilon(4S)$  dataset,
- Includes  $\Lambda_c^+$  reconstructed as  $pK^-\pi^+$ ,  $pK_s^0$ , and  $\Lambda\pi^+$ .
- Subdecay  $\overline{B}{}^0 \rightarrow \Sigma^0_c (\Lambda^+_c \pi^-) \pi^+ \overline{\rho}$ 
  - Yield from simultaneous fit to
    - Beam-constrained  $\overline{B}^0$  mass,  $M_{\rm bc}$ ,
    - $M(\Lambda_{c}^{+}\pi^{-}).$
  - Subset used in subsequent quasi-3-body analysis:  $|M(\Lambda_c^+\pi^-) - m_{PDG}(\Sigma_c^0)| < 14 \,\text{MeV}.$

Similarly for  $\overline{B}{}^0 \rightarrow \Sigma_c^{++}(\Lambda_c^+\pi^+)\pi^-\overline{\rho}$ .





#### EVIDENCE FOR A NEW BARYON STATE Belle: Phys.Rev.Lett. 130 (2023) 3, 031901 (arXiv:2206.08822 [hep-ex])

#### Coincident peaking features in

- $M(\Sigma_{c}^{0}(\Lambda_{c}^{+}\pi^{-})\pi^{+})$ , and
- $M(\Sigma_{c}^{++}(\Lambda_{c}^{+}\pi^{+})\pi^{-}).$

Not observed in  $\overline{B}{}^0$  and  $\Sigma_c^{++(0)}$  mass sidebands.

Provisionally assumed to be a new state,

• Tentatively named  $\Lambda_c(2910)^+$ .

Fits to combined  $M\left(\Sigma_c^{++(0)}(\Lambda_c^+\pi^{-(+)})\pi^{+(-)}\right)$ 

- 4.2 $\sigma$  total yield significance,
- $M(\Lambda_c(2910)^+) = 2913.8 \pm 5.6 \pm 3.8 \,\mathrm{MeV}_{\odot}$
- $\Gamma(\Lambda_c(2910)^+) = 51.8 \pm 20.0 \pm 18.8 \,\mathrm{MeV}.$



MORE  $\Omega_c(X)^0$  STATES AT LHCb: LHCB-PAPER-2022-043 NEW BARYON RESONANCES

More  $\Omega_c(X)^0$  states! LHCb: LHCb-PAPER-2022-043 (ArXiv:2302.04733 [HEP-EX])



Updates Phys.Rev.Lett. 118 (2017) 18, 182001, the analysis of the

 $M(\Xi_{c}^{+}K^{-})$  spectrum, to the full Run 1 + Run 2 data set.

• Approximately 5× more signal due to Run 2  $\sqrt{s}$  and trigger.

Two additional observed peaks!





MASSES AND WIDTHS OF  $\Omega_c(X)^0$  STATES LHCb: <u>LHCb-PAPER-2022-043</u> (ARXIV:2302.04733 [HEP-EX])

Masses and widths measured with improved precision over previous measurements.

			Also	
Resonance	<i>m</i> (MeV)	Г (MeV)	seen	
$\Omega_{c}(3000)^{0}$	$3000.44 \pm 0.07 \ ^{+0.07}_{-0.13} \pm 0.23$	$3.83 \pm 0.23  {}^{+1.59}_{-0.29}$	[1],[2]	
$\Omega_{c}(3050)^{0}$	$3050.18 \pm 0.04  {}^{+0.06}_{-0.07} \pm 0.23$	$0.67 \pm 0.17  {}^{+0.64}_{-0.72}$	[1],[2]	
		< 1.8 MeV, 95% C.L.		
$\Omega_{c}(3065)^{0}$	$3065.63 \pm 0.06 \ ^{+0.06}_{-0.06} \pm 0.23$	$3.79 \pm 0.20  {}^{+0.38}_{-0.47}$	[1],[2]	
$\Omega_{c}(3090)^{0}$	$3090.16 \pm 0.11 \stackrel{+0.06}{_{-0.10}} \pm 0.23$	$8.48 \pm 0.44  {}^{+0.61}_{-1.62}$	[1],[2]	
$\Omega_{c}(3119)^{0}$	$3118.98 \pm 0.12  {+0.09 \atop -0.23} \pm 0.23$	$0.60 \pm 0.63  {}^{+0.90}_{-1.05}$		
	0.20	< 2.5 MeV, 95% C.L.		
$\Omega_{c}(3185)^{0}$	$3185.1 \pm 1.7  {}^{+7.4}_{-0.9} \pm 0.2$	$50\pm7~^{+10}_{-20}$		
$\Omega_{c}(3327)^{0}$	$3327.1 \pm 1.2  {+0.1 \atop -1.3} \pm 0.2$	$20 \pm 5  {}^{+\overline{13}}_{-1}$		
[1] Belle $e^+e^- \rightarrow \Xi_c^+K^-X$ , Phys.Rev.D 97 (2018) 5, 051102				
[2] LHCb $\Omega_b^- \to \Xi_c^+ K^- \pi^+$ , Phys.Rev.D 104 (2021) 9, L091102				

# Particle properties

## BELLE

MASS AND WIDTH OF Λ<sub>c</sub>(2625)<sup>+</sup> Belle: Phys.Rev.D 107 (2023) 3, 032008 (<u>arXiv:2212.04062 [hep-ex]</u>)

Analysis of 
$$\Lambda_c(2625)^+ \rightarrow \Lambda_c^+ \pi^- \pi^+$$
  
in full Belle dataset

- $\mathcal{L}_{int} = 980 \, \text{fb}^{-1}$ ,
- $\Lambda_c^+$  subdecay to  $p^+ K^- \pi^+$ ,
- Approximately 30k signal decays.



Mass and width of  $\Lambda_c(2625)^+$  measured with  $M(\Lambda_c^+\pi^-\pi^+)$  distribution

- Kinematic fit of the decay with  $\Lambda_c^+$  constrained to world average,
- Width consistent with experimental resolution.

	$m(\Lambda_c(2625)^+) - m(\Lambda_c^+)$ (MeV)	$\Gamma(\Lambda_{c}(2625)^{+}) (MeV)$
This result	$341.518 \pm 0.006 \pm 0.049$	< 0.52
PDG 2022 <sup>1</sup>	$341.65 \pm 0.13$	< 0.97

<sup>1</sup>PDG 2022 WA dominated by CDF: Phys.Rev.D 84 (2011) 012003

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PARTICLE PROPERTIES  $\Lambda_c(2625)^+$  PROPERTIES AT BELLE: PRD 107 (2023) 3, 032008



AMPLITUDE ANALYSIS OF  $\Lambda_c(2625)^+ \rightarrow \Lambda_c^+ \pi^- \pi^+$ Belle: Phys.Rev.D 107 (2023) 3, 032008 (arXiv:2212.04062 [hep-ex])



 $A_{c}(2625)^{+}$  PROPERTIES AT BELLE: PRD 107 (2023) 3, 032008 PARTICLE PROPERTIES

SIDEBAND ANALYSIS OF  $\Lambda_c(2625)^+ \rightarrow \Lambda_c^+ \pi^- \pi^+$ Belle: Phys.Rev.D 107 (2023) 3, 032008 (arXiv:2212.04062 [hep-ex])

Non- $\Lambda_c(2625)^+$  contributions to  $\Sigma_c^{++(0)}\pi^{-(+)}$ estimated by analysis of  $M(\Lambda_c^+\pi^-\pi^+)$  sidebands,

- 6 sidebands (3 on each side),
- Linear extrapolation to  $\Lambda_c(2625)^+$  signal region.

$$\frac{\mathcal{B}(\Lambda_{c}(2625)^{+} \to \Sigma_{c}^{++(0)} \pi^{-(+)})}{\mathcal{B}(\Lambda_{c}(2625)^{+} \to \Lambda_{c}^{+} \pi^{-} \pi^{+})} = \frac{N_{sig}(\Sigma_{c}^{++(0)}) - N_{bkg}(\Sigma_{c}^{++(0)})}{N_{sig}(\Lambda_{c}(2625)^{+})}$$

$$\frac{\mathcal{B}(\Lambda_c(2625)^+ \to \Sigma_c^{++}\pi^-)}{\mathcal{B}(\Lambda_c(2625)^+ \to \Lambda_c^+\pi^-\pi^+)} = (5.13 \pm 0.26 \pm 0.32)\%$$

$$\frac{\mathcal{B}(\Lambda_c(2625)^+ \to \Sigma_c^0 \pi^+)}{\mathcal{B}(\Lambda_c(2625)^+ \to \Lambda_c^+ \pi^- \pi^+)} = (5.19 \pm 0.23 \pm 0.40)\%$$









### SPIN AND PARITY OF $D^*_{(s)}$ BESIII: ARXIV:2305.14631 [HEP-EX]

Helicity amplitude analysis in 3.19 fb<sup>-1</sup> of  $e^+e^-$  at  $\sqrt{s} = 4.178$  GeV

$$e^{+}e^{-} \rightarrow \gamma^{*} \rightarrow D^{*0}\overline{D}^{0}, D^{*0} \rightarrow D^{0}\pi^{0}$$

$$e^{+}e^{-} \rightarrow \gamma^{*} \rightarrow D^{*+}D^{-}, D^{*+} \rightarrow D^{+}\pi^{0}$$

$$e^{+}e^{-} \rightarrow \gamma^{*} \rightarrow D^{*+}_{s}D^{-}_{s}, D^{*+}_{s} \rightarrow D^{+}_{s}\gamma$$

$$e^{-}D^{*}_{s} \rightarrow D^{*+}_{s}D^{-}_{s} \rightarrow D^{*+}_{s}D^{+}_{s}\gamma$$

$$P^{*+}_{s} \rightarrow D^{*+}_{s}D^{*+}_{s} \rightarrow D^{*+}_{s}D^{$$

Likelihoods for each natural  $J^P$  in  $\{1^-, 2^+, 3^-\}$  constructed from helicity amplitudes,

$$\mathcal{L}^{J^{P}} = \prod_{i=1}^{N_{\text{evis}}} \frac{1}{\mathcal{C}} \mathcal{W}^{J^{P}}(\theta_{0}^{i}, \theta_{1}^{i}, \phi_{1}^{i}, m_{12}) = \prod_{i=1}^{N_{\text{evis}}} \frac{1}{\mathcal{C}} \overline{\sum_{m, \lambda_{i}}} \left| \mathcal{A}(m, \lambda_{1}, \theta_{0}^{i}, \theta_{1}^{i}, \phi_{1}^{i}, m_{12}) \right|^{2}$$

where *m* is the helicity of  $\gamma^*$  and  $\lambda_1$  the helicity of the  $\gamma$  or  $\pi^0$  from the  $D^*$ .

**PARTICLE PROPERTIES** 

 $J^P$  of  $D^*_{(s)}$  at BESIII: ARXIV:2305.14631 [HEP-EX]

## ₩S

## SPIN AND PARITY OF $D^*_{(s)}$ BESIII: arXiv:2305.14631 [hep-ex]



For each of  $D_s^{*+}$ ,  $D^{*0}$ , and  $D^{*+}$ ,  $J^P = 1^-$  fits well.

- $J^P = 2^+$  and  $3^-$  tested as 'null hypotheses'
  - Likelihood compared to that of a linear combination of  $J^P = 1^-$  and the 'null'  $J^P$ .

Process	Hypothesis	$\Delta(-2\ln \mathcal{L})$
*	1 <sup>-</sup> over 2 <sup>+</sup>	1102
$D_s \cdot D_s$	1- over 3-	2104
0م 0*0	1 <sup>-</sup> over 2 <sup>+</sup>	12134
$D^{-1}D^{+1}$	1 <sup>-</sup> over 3 <sup>-</sup>	12096
-0++0-	1 <sup>-</sup> over 2 <sup>+</sup>	11308
	1- over 3-	11222

In all cases, the bare 'null'  $J^P = 2^+$  and  $3^-$  disfavored by  $> 10\sigma$ .

# Strong and EM decays

OBS. OF  $D^{*0} \rightarrow D^0 e^+ e^-$  AT BESIII: PRD 104 (2021) 11, 112012 STRONG AND EM DECAYS



OBSERVATION AND BR OF  $D^{*0} \rightarrow D^0 e^+ e^-$ BESIII: PHYS.REV.D 104 (2021) 11, 112012 (ARXIV:2111.06598 [HEP-EX])

Analysis in  $e^+e^- \rightarrow D^{*0}\overline{D}^{*0}$  events from  $3.19 \,\mathrm{fb}^{-1}$  at  $\sqrt{s} = 4.178 \,\mathrm{GeV}$ .

 $D^0$  reconstructed in decay modes  $K^-\pi^+$ ,  $K^{-}\pi^{+}\pi^{0}$ , and  $K^{-}\pi^{+}\pi^{-}\pi^{+}$ .

Backgrounds from  $D^{*0} \rightarrow D^0 \gamma$  via  $\gamma \rightarrow e^+ e^$ material conversions suppressed by  $e^+e^$ vertex cut.

Clear  $D^{*0} \rightarrow D^0 e^+ e^-$  peak in each  $D^0$  mode,

• 13.2 $\sigma$  total statistical significance.

Branching ratio measured relative to  $D^{*0} \rightarrow D^0 \gamma$ :  $rac{\mathcal{B}(D^{*0} \to D^0 e^+ e^-)}{\mathcal{B}(D^{*0} \to D^0 \gamma)} = (11.08 \pm 0.76 \pm 0.49) imes 10^{-3}$ 



STRONG AND EM DECAYS BF OF  $D_s^{*+} \rightarrow D_s^+ \pi^0$  at BESIII: PRD 107 (2023) 3, 032011

BRANCHING FRACTION OF  $D_s^{*+} \rightarrow D_s^+ \pi^0$ BESIII: Phys.Rev.D 107 (2023) 3, 032011 (ArXiv:2212.13361 [hep-ex])

Analysis of  $e^+e^- \rightarrow D_s^{*+}D_s^-$  events from 7.33 fb<sup>-1</sup> at  $\sqrt{s} = 4.128$  to 4.226 GeV.

Reconstruct both  $D_s^+$  and  $D_s^-$  in each event,

• In decays to either  $K^{\pm}K^{\mp}\pi^{\pm}$  or  $K^0_{\rm s}K^{\pm}$ 

Spectrum of squared missing mass,

$$M_{
m miss}^2 \equiv (\sqrt{s} - E_{D_s^+} - E_{D_s^-})^2 - |ec{p}_{D_s^+} + ec{p}_{D_s^-}|^2,$$

• Peaks at 0 for 
$$D_s^{*\pm} \rightarrow D_s^{\pm} \gamma$$
,

• Peaks at  $m^2_{\pi^0}$  for  $D^{*\pm}_s o D^\pm_s \pi^0.$ 

Cut-and-count methodology,

- Partition into  $\gamma$  and  $\pi^0$  regions,
- Unfold with an efficiency matrix.

 $rac{D_s^{++} o D_s^{+} \pi^0}{D_s^{++} o D_s^{+} \gamma} = (6.16 \pm 0.43 \pm 0.18)\%$ 









# Doubly-charmed baryons

## $\Xi_{\rm cc}$ at LHCb so far







#### With properties:

- Mass  $m = 3621.55 \pm 0.23 \pm 0.30$  MeV (JHEP 02 (2020) 049)
- Lifetime  $\tau = 0.256^{+0.024}_{-0.022} \pm 0.014 \, \text{ps}$  (Phys.Rev.Lett. 121 (2018) 5, 052002)

• 
$$\left(\frac{\sigma(\Xi_{cc}^{++})\mathcal{B}(\Xi_{cc}^{++}\to\Lambda_{c}^{+}K^{-}\pi^{+}\pi^{+})}{\sigma\Lambda_{c}^{+}}\right)_{\text{fid}} = (2.22 \pm 0.27 \pm 0.29) \times 10^{-4}$$
  
(Chin.Phys.C 44 (2020) 2, 022001)

#### OBSERVATION OF $\Xi_{cc}^{++} \rightarrow \Xi_{c}^{\prime+}\pi^{+}$ LHCb:JHEP 05 (2022) 038 (ARXIV:2202.05648 [HEP-EX])



Search in the  $\Xi_c^+\pi^+$  mass spectrum (with  $\Xi_c^+ \rightarrow \rho K^-\pi^+$ ),

- Signature of  $\Xi_{cc}^{++} \rightarrow \Xi_{c}^{\prime+} \pi^{+}, \ \Xi_{c}^{\prime+} \rightarrow \Xi_{c}^{+} \gamma$  with unreconstructed  $\gamma$ ,
- Appears as a peaking structure at a lower mass than the previously observed  $\Xi_{cc}^{++} \rightarrow \Xi_{c}^{+}\pi^{+}$  signal.

Analyzed in two statistically independent subsets



10 SEARCH FOR 2 66 AT EFFECT, STIEL 12 (2021)



Search with a two-tiered selection.

Most restrictive: default trigger set,

- Selection and trigger requirements matched to  $\Xi_{cc}^{++} \rightarrow \Xi_{c}^{+}\pi^{+},$
- Best precision/upper limits of production × branching ratio.





Superset: *extended trigger* set.

• More inclusive set of triggers.

Best chance of detecting a signal.

Full analysis, including decision tree of how to present results in all scenarios, fully defined before unblinding. **DOUBLY-CHARMED BARYONS** 

CONTINUING SEARCH FOR  $\Xi_{CC}^+$  AT LHCb:, JHEP 12 (2021) 107



COMBINATION OF  $\Xi_c^+\pi^-\pi^+$  and  $\Lambda_c^+K^-\pi^+$ LHCb:JHEP 12 (2021) 107 (arXiv:2109.07292 [hep-ex])

No evidence for 
$$\Xi_{cc}^+ \to \Xi_c^+ \pi^- \pi^+$$
.

Determination of mass-dependent upper limit for

 $\boldsymbol{R} \equiv \frac{\sigma(\boldsymbol{\Xi}_{cc}^{+}) \times \mathcal{B}(\boldsymbol{\Xi}_{cc}^{+} \rightarrow \boldsymbol{\Xi}_{c}^{+} \pi^{-} \pi^{+})}{\sigma(\boldsymbol{\Xi}_{cc}^{++}) \times \mathcal{B}(\boldsymbol{\Xi}_{cc}^{++} \rightarrow \boldsymbol{\Xi}_{c}^{+} \pi^{+})}$ 

*R* estimated for several hypothetical lifetimes.





Statistical combination with  $\Xi_{cc}^+ \rightarrow \Lambda_c^+ K^- \pi^+$  search

Sci.China Phys.Mech.Astron. 63 (2020) 2, 221062

Joint fit to the mass spectra.

No evidence for  $\Xi_{cc}^+$ 

 Largest deviation: 2.9σ after systematics and look-elsewhere. SUMMARY

#### WRAP-UP AND ACKNOWLEDGMENTS

Our experimental understanding of charm spectroscopy continues to benefit from complementary experiments,

- Direct production in both  $e^+e^-$  and pp collisions,
- Production in decay of large *b*-hadron samples.

The continuously increasing sample sizes are furthering both the precision and comprehensiveness of our knowledge.

Large-sample flavor experiments continue to discover new and interesting behaviors of charmed hadrons.



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